

# Signal Significance in Particle Physics

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## Abstract

The concept of the “statistical significance” of an observation, and how it is used in particle physics experiments is reviewed. More properly known as a “p-value,” the statistical foundations for this concept are reviewed from a frequentistic perspective. The discovery of the top quark at the Fermilab Tevatron Collider and a more recent analysis of data recorded at Fermilab are used to illustrate practical applications of these concepts.<sup>1</sup>

## 1. What Particle Physicists Mean by Significance

When one of your colleagues approaches you and declares that she has made a “significant” observation, intuitively that means that she has observed some phenomena whose interpretation allows her to eliminate or falsify one or more hypotheses, and usually support one or a small number of alternative hypotheses. We furthermore expect that the observation has sufficient statistical power that we expect that additional observations are unlikely to change these conclusions. Scientists have therefore attempted to identify a consistent statistical framework in which we can quantify this concept of “significance.”

In particle physics, this concept of statistical significance has not been employed consistently in the most important discoveries made over the last quarter century. Examples of the major discoveries made over an approximately 10 year period between the late 1970’s and the late 1980’s illustrate this point.

Let us consider first the discovery of the  $\Upsilon$  meson (and the  $b$  quark) in 1977 by L. Lederman and colleagues [1]. This was made through the observation of  $\mu^+\mu^-$  final states in high-energy proton-nucleus collisions at Fermilab, where a large resonant signal was observed on top of a steeply falling background of dimuon candidates. The experimenters estimated that they observed a signal of approximately 770 events on top of a non-resonant background of 350 candidates. They characterized the signal as “significant” but made no attempt to quantify or explain exactly what they meant.

The discovery of the  $W^-$  boson at CERN in 1983 by the UA1 collaboration [2] was made by observing 6 events produced in proton-antiproton collisions where a high energy electron or antielectron was observed in coincidence with a signature for a recoiling energetic neutrino. The collaboration estimated the background to these 6 events as being “negligible” and claimed discovery of the expected charged weak intermediate vector boson. This observation was subsequently confirmed by the UA2 collaboration.

The discovery of  $B$  mesons in 1983 by the CLEO collaboration [3] was performed by carefully reconstructing a variety of different decay modes and illustrating an invariant mass peak at  $5.4 \text{ GeV}/c^2$ . The collaboration observed a total event rate of 17 events on a background of between 4 and 7 events. They

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claimed definitive observation of a new particle, but made no statement that quantified the statistical power of the observation.

As a final example, I note the discovery of  $B^0$  meson flavour mixing in 1987 by the ARGUS collaboration [4]. The experimenters observed  $24.8 \pm 7.6 \pm 3.8$  unexpected same-sign dilepton events versus a total of  $25.2 \pm 5.0 \pm 3.8$  opposite-sign dilepton candidates. They characterized this as a “3  $\sigma$ ” observation, namely, that the probability that the observed number of same-sign dilepton events could have been as great or greater than the observed value was equivalent to the probability of a Gaussian statistic being observed at least 3 standard deviations from its expected mean (a probability of  $1.35 \times 10^{-3}$ ).

This brief review illustrates that quantifying the statistical significance or power in seminal particle physics measurements is not uniformly done. It also illustrates that in at least the one case in which it was done, the significance was defined as the probability<sup>2</sup> of the “null hypothesis” having been responsible for the observation.

In this paper, I will first review briefly the formal concept of “statistical significance.” I will then discuss several examples that illustrate the use of this concept in particle physics. I do not have the opportunity to review all of the techniques that have been in recent use, but refer the reader to other articles in these proceedings (for example, the review of the  $CL_S$  method by A. Read).

## 2. Formal Definitions of Significance

### 2.1 The Frequentists Perspective

The concept of statistical significance is formally introduced in the context of hypothesis testing [5]. Suppose that we have two hypotheses,  $H_0$  and  $H_1$ , and a measurement whose value is a test statistic  $X$  that, as a random variable, provides some discrimination between these two hypotheses. Let  $f_0(X)$  and  $f_1(X)$  represent the probability distribution functions for  $X$  associated with the two hypotheses.

Prior to making a measurement of  $X$ , we would identify a “critical region,”  $w$ , such that we would select hypothesis  $H_1$  if  $X \in w$  and  $H_0$  otherwise. We now have four possible outcomes when we make a measurement of  $X$ . If  $X \in w$  and the hypothesis  $H_1$  is true, then we have selected the correct hypothesis. If  $X \in w$  and  $H_0$  is true, then we have incorrectly concluded that  $H_1$  is true. This is known as a mistake of the first kind, and the probability for this decision is

$$\int_{X \in w} f_0(X) dX = \alpha. \quad (1)$$

The probability  $\alpha$  is known as the “significance” of the test.

We have two other possibilities. The first is if we measure  $X \notin w$  when  $H_0$  is true. In this case, we would have made the correct inference. Finally, we have the case where  $X \notin w$  and  $H_1$  is true. This is known as a mistake of the second kind, and the probability for that decision is

$$\int_{X \notin w} f_1(X) dX = \beta. \quad (2)$$

The probability  $1 - \beta$  is known as the “power” of the test. The situation is illustrated in Fig. 1a). The significance  $\alpha$  is therefore a measure of the ability of a test to avoid mistakes of the first kind, whereas the power  $1 - \beta$  measures the ability of a test to avoid mistakes of the second kind. In defining an “optimum” test, one would like to choose  $X$  and the region  $w$  such that  $\alpha$  and  $\beta$  are as small as possible.

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<sup>2</sup>Unless otherwise noted, “probability” in this article refers to the frequentist definition of this concept.

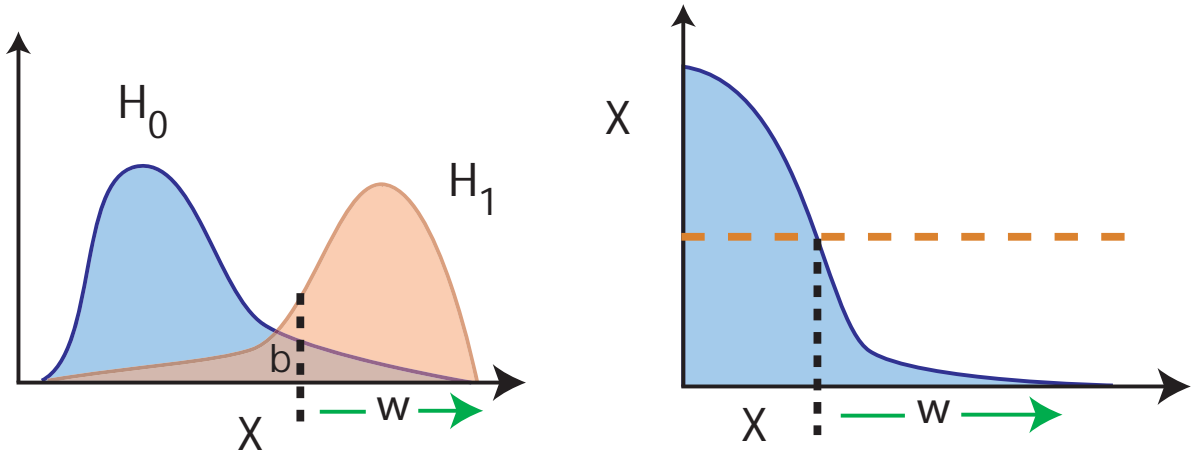


Fig. 1: A schematic of the hypothesis test described in the text is shown in a). The plot shows the probability densities for  $X$  under the two hypotheses  $H_0$  and  $H_1$  and one possible choice for the region  $w$ . The use of the Neyman-Pearson Theorem is illustrated in b), where the ratio  $I_N$  is plotted as a function of  $X$ .

## 2.2 Significance in Particle Physicists– The P-Value

The statistical definition of significance is made in the context of choosing between two hypotheses. However, the use of significance in particle physics discoveries is in a different context. The typical case is that an experiment makes a measurement of the test statistic  $X$ , say  $X_0$ . Furthermore, using the same notation as before, we assume  $X$  has a probability density  $f_0(X)$  if the hypothesis  $H_0$  is true. We further assume that we can categorize observations of  $X$  into those that are more and less consistent with  $H_0$  (for the sake of discussion, I will assume that values of  $X$  greater than  $X_0$  are less likely given  $H_0$ ). A measure of the inconsistency of the observed value  $X_0$  with the hypothesis  $H_0$  is then the probability

$$\int_{X > X_0} f_0(X) dX. \quad (3)$$

This is formally identical to the definition of the significance in Eq. 1 if we now define the critical region to be  $w = \{X | X > X_0\}$ , i.e. the probability of observing a value of  $X$  equal to or greater than our observation. This probability is formally known as the “p-value” of the observation, a convention that the Particle Data Group now has adopted [6]. The advantage of using the formal term for this quantity is that it avoids confusion with the concept of significance defined in hypothesis tests, where the region  $w$  is defined *a priori*, i.e. before the measurement is made.

The p-value for a given measurement and a specific hypothesis has a number of features. First, it only depends on the measurement and the probability density for the hypothesis. It is not a hypothesis test. It only provides a measure of the consistency of the hypothesis and the measurement. In that sense, it is most often quoted when one has made a measurement that appears to be inconsistent with a single hypothesis. A very small p-value is then used to support the inference that the specific hypothesis should be rejected.

Referring then to the example of the discovery of  $B^0$  mixing given in Section 1, we can now say that the p-value for the observation for the non-mixing hypothesis was  $1.35 \times 10^{-3}$ . From a frequentist perspective, if one rejected the non-mixing hypothesis at this p-value and it was always true, then one would expect to be wrong (*i.e.*, reject the correct hypothesis) on average 1 out of every 740 times.

### 2.3 A Few More Comments on Hypothesis Testing

Although the literature uses the p-value of an observation as a measure of its statistical significance, the concept of hypothesis testing is an important one in particle physics. One sees it most often used in the context where one is designing or proposing an experiment and wishes to characterize the experiment's ability to distinguish existing and known physics phenomena (such as that predicted by the Standard Model) from possible new physics [7].

In those cases, a crucial aspect of the experiment design is the selection of the optimal statistic  $X$  and the optimal critical region. For a specific measurement, such as the observation of a process above some expected background rate, the choice of  $X$  will depend on the measurement and the ingenuity of the experimenter. One would like to identify test statistics that have quite different probability density functions for the hypotheses you wish to distinguish in order to be able to define a critical region with the smallest possible  $\alpha$  and  $\beta$ . At the same time, the decision should be informed by the effect of any systematic uncertainties that may degrade the separation between two hypotheses. To that extent, one often attempts to identify statistics that are not expected to be affected by systematic uncertainties.

### 2.4 Neyman-Pearson Theorem

Hypothesis testing has an extensive literature, but relatively few general results have been identified that can guide our judgement. One result, known as the Neyman-Pearson Theorem, is surprisingly useful, and it is worth reviewing here for the insight it provides.

Suppose we have two hypotheses,  $H_0$  and  $H_1$ , and we have defined a test statistic  $X$ . Then for a given significance  $\alpha$ , we can define the region  $w$  which gives us optimal  $\beta$  (i.e., the smallest value of  $\beta$  and therefore the greatest statistical power) by choosing  $w$  as follows. We first form the ratio of probability density functions for the two hypotheses

$$I_N(X) \equiv \frac{f_0(X)}{f_1(X)}. \quad (4)$$

The Neyman-Pearson Theorem then concludes that the optimal region  $w$  is the one over which  $I_N(X)$  is maximal, namely that we find the value  $c_\alpha$  such that when

$$w = \{X | I_N(X) < c_\alpha\}, \quad (5)$$

the probability of observing  $X \in w$  is

$$\int_{X \in w} f_0(X) dX = \alpha. \quad (6)$$

This construction is illustrated in Fig. 1b).

The Neyman-Pearson test has one significant limitation—it is only valid for what are known as “simple hypotheses,” or hypotheses where there are no unknown parameters that would be estimated from the data. In addition, since it is only applicable when comparing two hypotheses, it cannot be employed in cases where you have multiple alternative hypotheses to consider. However, despite these limitations, this theorem gives us considerable insight into the definition of the critical region. For example, we can relate the ratio of probabilities to the ratio of likelihoods of the two hypotheses:

$$I_N(X) \equiv \frac{f_0(X)}{f_1(X)} \sim \frac{L_0(X)}{L_1(X)}, \quad (7)$$

where  $L_i(X)$  are the likelihood functions defined for the two hypotheses  $i = 0$  and  $i = 1$ . This suggests that the likelihood ratio is one source of guidance for defining critical regions that have significant (though perhaps not optimal) power.

### 3. The Bayesian Perspective

Our considerations up to this point have been from a frequentist perspective, using the standard definition of a frequentist probability. In calculating a p-value for a measurement, one has to assume a hypothesis and then determine the probability (or probability density) for all possible outcomes of the measurement.

A Bayesian statistician does not consider data other than the single measurement. However, for each hypothesis, the Bayesian could define a credibility interval that reflects his or her degree-of-belief in each hypothesis, and the ratio of these credibility intervals—what is called the “Bayes discriminant factor”—becomes a measure of the relative confidence one has in the two hypothesis. Formally, this ratio is

$$\frac{P(H_0|X)}{P(H_1|X)} = \frac{L_0(X)}{L_1(X)} \frac{\pi_0(X)}{\pi_1(X)}, \quad (8)$$

where  $\pi_0(X)$  and  $\pi_1(X)$  are the prior probabilities associated with the two hypotheses.

This ratio can be used to reject one of the two hypotheses. The Bayesians would argue that there is no benefit in attempting to make anything but a relative statement about the degree-of-belief of the two hypotheses. Thus, there is no direct analogy to the p-value in this framework. The advantage of this perspective is that it avoids the need to understand the probability density of all possible outcomes for a given hypothesis. It also has the advantage that any inferences you draw are less sensitive to an outcome that has a low probability regardless of the hypothesis. In such cases, the Bayes discriminant factor still provides information, whereas the p-value is no longer very informative and could in fact be misleading.

The disadvantages with this Bayesian approach are, however, that one has to assume prior distributions for each hypothesis, and one is only allowed to make relative confidence statements about two hypotheses. For these reasons, one finds very limited use of the Bayes discriminant factor in particle physics.

### 4. P-Values and Experimental Design

The definition of significance in terms of a p-value for an observation immediately makes clear the importance of *a priori* decisions on the random variables one will measure and how one will define those observations that prefer one hypothesis over another. A carefully designed experiment will identify these and optimize their choice before any data is analyzed.

However, many particle physics experiments make unique measurements using general-purpose apparatus designed to study a large class of processes. Thus, it is difficult, and often impossible, to anticipate what one will observe and how. In fact, early studies of the data will often guide the experimenters to focus in specific features that appear unusual or unexpected. In this context, the evaluation of a p-value may prove very difficult.

A simple example illustrates this problem. Suppose one measures an invariant mass spectrum in a specific region, say  $[m_1, m_2]$ , and one observes a narrow enhancement in a small mass interval, say  $\Delta m$  wide, of  $N_o$  events above an expected background of  $N_b$  events. In this case, it would be natural to assume that the hypothesis we wish to test is the “null” hypothesis where we expect  $N_b$  events in this mass interval  $\Delta m$  and then determine the probability of observing at least  $N_o$  events. Assuming that the background rate

is well known (and so we can ignore its uncertainty), the p-value for this observation would be given by the Poisson probability for observing at least  $N_o$  events when the mean rate is  $N_b$ , or

$$\alpha = \sum_{n=N_o}^{\infty} \frac{\exp(-N_b) (N_b)^n}{n!}. \quad (9)$$

However, this probability does not take into account the fact that we are considering all possible choices of mass interval  $\Delta m$  in the region  $[m_1, m_2]$ .

A proper estimate of this p-value would then have to include the likelihood of observing at least  $N_o$  events in *any* possible interval  $\Delta m$ . This increases the p-value of the observation, and changes the possible inferences one can make. For example, a Monte Carlo calculation where  $\delta m$  is 1% of the interval,  $N_o = 8$  events and  $N_b = 100$  (i.e., the average number of events in any  $\delta m$  interval is one) gives a p-value that is 500 times larger than the result in Eq. 9.

#### 4.1 Blind Analyses

The prevalence of the p-value in making inferences rests on the assumption that it is possible to estimate the frequency of all observations of the test statistic, and that it is possible to identify the class of observations that are less consistent with a given hypothesis (the critical region in the language of hypothesis testing). This is inherently difficult in cases where one allows the definition of test statistic and critical region to depend on the actual experimental outcome itself. A tactic to eliminate such bias is the “blind analysis,” where one defines the critical region and the statistic without knowledge of the relevant data [8].

The ideal experiment is one in which the measurement and any calculation of its p-value does not have to be informed by the data itself. No choices with regard to selection of data, modifications in the test statistic or choice of critical region would then be allowed once data collection has started. This approach avoids the possibility of selecting, consciously or unconsciously, a critical region or test statistic that tend to favour or disfavour a given hypothesis *based on the data observed*.

A number of celebrated failures of inference in particle physics over the last half-century illustrate what happens when the experimenter allows the data to guide his or her choices in making inferences about data [8]. In all these cases, the quoted p-value has been assessed incorrectly because it has failed to take into account how the frequency of a given observation would be affected by making choices about the test statistic and critical region based on the actual distribution of the data itself.

#### 4.2 Use and Limitations of Blind Experiments

The simple example of “bump hunting” illustrates the fundamental problem in particle physics where one is searching for evidence of new phenomena; it is inherently difficult to identify *a priori* what class of observations one would expect to use in such a search. Besides the difficulty of defining in advance all possible means of separating “signal” from “background,” it is also difficult to limit access to data when one also has to verify that the instrumentation is working correctly and that any artefacts created by effects such as miscalibration and errors in bookkeeping are identified and mitigated. The experiment design also has to allow the experimenter access to the data to measure the rate of background events in the signal sample.

Despite these challenges, the elimination of certain biases that are otherwise difficult to control make a blind analysis an attractive approach given the benefits of being able to make straightforward estimates of p-values for the possible outcomes. This technique is reviewed in another contribution to these proceedings [8].

## 5. P-Values in a Counting Experiment

### 5.1 General Considerations

A common particle physics experiment involves the search for new phenomena by observing a unique class of events in particle interactions that cannot be described by background hypotheses. One usually can reduce this problem to that of a “counting experiment,” where one identifies a class of events using well-defined criteria, counts up the total number of observed events,  $N_o$ , and estimates the average rate of events,  $N_b$ , that come from the various background processes. One can then perform a straightforward estimate of the p-value of a given observation of  $N_o$  events, assuming that the probability density for the random variable  $N_o$  follows a Poisson distribution, *i.e.* the formula in Eq. 9.

There are several issues that even this simple problem has to address. First, one has to be sure that the criteria used to select the class of events was not in itself biased by how  $N_o$  varied as the criteria were modified. Here is where a blind analysis has its greatest benefit, since this bias is explicitly guarded against. Second, one has to take into account possible uncertainties in the estimate of the background rate  $N_b$ . It is beyond the scope of this article to discuss this issue, and the interested reader is referred to the growing literature on this topic [9] (a typical frequentist approach is to extend the ensemble of possible measurements to include those experiments with different values of  $N_b$  consistent with the knowledge of  $N_b$ ). Third, the careful experimenter has to make sure that all information relevant to the search is used in the measurement. It is at best inefficient and at worst misleading to ignore relevant data (for example, a possible channel in which the number of observed events can provide additional information on the process being studied).

As a concrete example of the calculation of a p-value for a typical counting experiment, I will summarize the techniques used by the CDF and D0 collaborations in their search for top quark production.

### 5.2 The Top Quark Search

The top quark was discovered by pair-production in proton-antiproton collisions at an energy of 1.8 TeV [10, 11]. The top quark decays predominantly via the process  $t \rightarrow Wb$ , with the  $W$  boson subsequently decaying either leptonically via  $W \rightarrow l\nu_l$  (where “ $l$ ” can be either an electron, muon or tau lepton) or hadronically via  $W \rightarrow q\bar{q}'$  (the quark final states are either  $u\bar{d}$  or  $c\bar{s}$ ). This results in three categories of possible final states with different topologies, efficiencies and background rates:

1. the lepton+jets channel, involving one high energy lepton, a neutrino and three or more jets from the hadronic decay of the  $W$  and the  $b$  quarks,
2. the dilepton channel, involving two high energy leptons, evidence for two neutrinos, and two or more jets from the  $b$  quarks, and
3. the hadronic channel, involving at least six jets.

In both experiments, one had to use additional criteria to improve the signal-to-noise ratios in the final candidate event samples. For CDF, the most effective way to do this was to require evidence that at least one of the jets arose from a  $b$  quark using two different “ $b$ -tagging” techniques. Thus, one could characterize the final states by the number of  $b$  tags, with the events with one or two  $b$  tags having increasing purity. For D0, the most effective way to reduce backgrounds was by imposing more stringent kinematic criteria (a topological selection) and using soft muon  $b$ -tagging.

The searches used data samples of increasing sensitivity. The first reported data came after the CDF and D0 collaborations had recorded 19.6 and 15.0 pb<sup>-1</sup>, respectively [12, 13]. At that time, the experiments had not completed analysis of the hadronic channels, which were expected to be dominated by background.

Final State	Observation (events)	Expected Background (events)	$\mathcal{B} \times \text{Efficiency}$	Expected Signal (events)
<b>CDF</b>				
Lepton + Jets (SVX $b$ -tags)	6	$2.3 \pm 0.3$	0.015	2.4
Lepton + Jets (Soft lepton $b$ -tags)	7	$3.1 \pm 0.3$	0.012	1.9
Dileptons	2	$0.6 \pm 0.3$	0.008	1.3
<b>D0</b>				
Lepton + Jets (Soft lepton $b$ -tags)	2	$0.6 \pm 0.2$	0.009	1.0
Lepton + Jets (Topology)	4	$1.8 \pm 0.9$	0.026	2.8
Dileptons	1	$0.8 \pm 0.1$	0.007	0.7

Table 1: The observed number of top quark candidates, the expected background rate, the overall branching ratio times efficiency for the channel, and the expected number of signal events assuming a top quark with a mass of  $160 \text{ GeV}/c^2$  for each final state.

Final State	P-Value
<b>CDF</b>	
Lepton + Jets (SVX $b$ -tagging)	0.032
Lepton + Jets (Soft lepton tagging)	0.038
Dileptons	0.012
Combined	0.0026
<b>D0</b>	
Combined	0.072

Table 2: The p-values determined for the observed event rates assuming the Standard Model background processes by the CDF and D0 collaborations. The D0 collaboration only reported a p-value for the observation of 7 candidate events with an expected background of  $3.2 \pm 1.1$  events.

The results of these analyses are summarized in Table 1, where we list the number of observed events, the estimated background rates, and the branching ratio times efficiency of observing a  $t\bar{t}$  decay in each mode.

The collaborations evaluated the statistical significance of their data by using a Monte Carlo calculation to estimate the frequency that the expected background processes would create a combined signal that was at least as large as that observed. The Monte Carlo calculation created an ensemble of experiments that modelled the possible observations in all channels assuming the Standard Model background hypothesis. For a given channel, the estimated background rate was used as the mean of a Poisson distribution of observed events. In order to account for uncertainties in the background rate, the mean value used to generate a new member of the ensemble was obtained by sampling a Gaussian distribution with the mean and width of the estimated background rate [9]. The results of these p-value calculations are summarized in Table 2. The collaborations concluded that the individual observations did not provide sufficient evidence to exclude the background hypothesis.

The collaborations proceeded to determine how likely their set of observations were assuming the background hypothesis by identifying a statistic that combined the observations in the individual channels. In the case of a counting experiment involving several channels, the maximum likelihood estimate of the rate of the process is simply the sum of the event rates in each channel. Thus, the natural statistic to evaluate



the combined significance of the observations was the observed sum of events in all channels. However, the CDF collaboration noted that the most sensitive measure of the cross section was not the total number of observed events in their sample, but the total number of observed  $b$ -tags (since there was a much larger probability of observing two  $b$ -tags in a signal event than in an event from a background process). Thus, CDF chose as its statistic the sum of the number of  $b$ -tags in the lepton+jet events combined and the number of dilepton events. Since the D0 data relied less on  $b$ -tagging, the collaboration chose to use the total number of observed events.

The calculation of the p-value of the observation assuming the background hypothesis was performed by a Monte Carlo procedure that effectively created a set of “pseudo-experiments.” In each pseudo-experiment, the number of  $b$ -tags and dilepton events from the different background sources was drawn from a Poisson distribution that had as its mean value the estimated background rate for the process. The uncertainty in the various background components was taken into account as described above, as was the correlation in the different background sources. This correlation arose from the fact that a number of background sources contributed both types of  $b$ -tags, whereas others did not. In effect, this increased the frequency of observing a larger number of  $b$ -tags (since now the fluctuations in the two components were correlated).

The resulting p-values are summarized in Table 2. One sees that the single most significant p-value was  $2.6 \times 10^{-3}$ . If one had not taken into account the correlations between the background sources, the combined p-value would have been  $1.6 \times 10^{-3}$ , or a factor of almost two smaller. Alternatively, the combined p-value determined by just counting events would have been approximately  $10^{-2}$ . This demonstrates the sensitivity of a p-value calculation to the approximations used to determine it. Given all this information, both experiments concluded that the observations were not sufficiently compelling statistically to exclude the background hypothesis.

### 5.3 Significance Required for Discovery

In the search for the top quark, the CDF and D0 collaborations argued that observations with p-values of order  $10^{-3}$  were not sufficiently significant to be used to claim discovery of a new phenomenon. Although this is clearly a matter of opinion, it is roughly consistent with the practice in the field, where typically the “ $5\sigma$ ” standard is used as rough rule of thumb to define the sensitivity necessary for discovery. This corresponds to a p-value equivalent to between  $5.7 \times 10^{-7}$  and  $2.8 \times 10^{-7}$ , depending on whether you are searching for a deviation from a mean or a one-sided fluctuation from the mean.

As a concrete example, the two Tevatron collaborations used an identical analysis procedure when approximately a factor of two more data had been recorded by both experiments. The resulting p-values of the CDF and D0 observations assuming the background hypothesis were  $1 \times 10^{-6}$  and  $2 \times 10^{-6}$ , respectively [10, 11]. Both experiments concluded that the background hypothesis could be excluded and claimed observation of top quark pair production.

## 6. P-Values for Continuous Test Statistics

High-energy physics measurements often examine statistical variables that are continuous in nature. In fact, to identify a sample of events enriched in the signal process, one often imposes selection requirements on such continuous variables. Often, it is important to take into account the entire distribution of a given variable for a set of events, and not just whether the events lie in a given range of values.

The general problem can be posed in the following way. Suppose we have a set of event data each characterized by a set of statistics  $\vec{X}_i$ , where  $i = 1$  to  $N$ . In addition, one has a hypothesis to test that predicts

the distribution of  $\vec{X}$ , say  $f(\vec{X}; \vec{\alpha})$ , where we assume this function to be normalized to unity between  $X_{min}$  and  $X_{max}$ , the minimum and maximum values of  $X$ , and  $\vec{\alpha}$  is a set of parameters that are either known or estimated directly from the data. Then the general problem is to define a statistic that gives a measure of the consistency of the distribution of data with the distribution given by the hypothesis.

### 6.1 Possible Tools

The most widely used such statistic in the 1-dimensional case is a form of a “runs test,” which compares the predicted cumulative distribution

$$g(X) = \int_{X_{min}}^X f(X') dX' \quad (10)$$

with the observed cumulative distribution  $h(X)$ . The most common test is the Kolmogorov-Smirnov (K-S) test [14], which makes this comparison by first finding the K-S distance

$$\delta = \max \{|g(X) - h(X)|, X \in (X_{min}, X_{max})\}, \quad (11)$$

namely the largest difference between the two cumulative distributions. This test statistic has a characteristic distribution that can be calculated analytically to provide one with a p-value, specifically the probability that one would observe a value of this test statistic as large as or larger than the observed value.

The K-S test gives a distribution-free measure of the consistency of a 1-dimensional continuous variable and is often used in the particle physics literature. Although there are a number of other tests that could be used in this case, all with similar properties [15], the K-S test has become a reference standard to employ.

### 6.2 Extension to Higher Dimensions

The K-S test (and other runs tests) are in principle limited to 1-dimensional distributions, but there are extensions to the case of several dimensions, though with a number of restrictions. The extension requires one to assume that the probability distribution predicted by the hypothesis can be factorized, so that

$$f(\vec{X}) = f_1(X_1)f_2(X_2) \cdots f_n(X_n), \quad (12)$$

where  $n$  is the number of continuous variables being compared. This in effect requires each of the variables to be uncorrelated, a strong assumption and one that has to be verified in practice. With this assumption, however, one can then define a set of independent statistics  $\delta_i$ ,  $i = 1$  to  $n$ , and the associated p-value for each observed K-S statistic  $p_i$ . Then one can combine these independent p-values into a single measure of significance.

### 6.3 Example: CDF “Superjets”

A concrete example of this technique is a recent analysis of hadron collider data performed by the CDF collaboration. A study was performed of events that were consistent with the production of one or more hadronic jets and a  $W$  boson decaying to a lepton-neutrino pair. The collaboration defined a subsample of these events where at least one jet was identified as a “superjet”, namely a  $b$ -quark candidate jet with both the presence of a secondary vertex in the jet displaced from the interaction vertex and the presence of a second lepton associated with the jet consistent with coming from the semileptonic decay of a  $b$  hadron [16].

The collaboration found 13 such events in the 1992-96 Tevatron Collider data, where they estimated that they would have expected  $4.4 \pm 0.6$  events from Standard Model background sources. This observation has a p-value of 0.001, treating it as a counting experiment and using the techniques introduced above. The authors then proceeded to examine nine separate kinematic variables that had distributions that were predicted to be largely uncorrelated, but that might distinguish between the Standard Model backgrounds and a variety of exotic sources of events. A typical example of such a comparison is given in Figure 2, where the observed distribution of the lepton pseudorapidity ( $\eta \equiv -\ln \tan(\theta/2)$ , where  $\theta$  is the angle of the lepton relative to the incoming proton beam axis) is compared with the predicted  $\eta$  distribution.<sup>3</sup> The plots on the right-hand side are the distributions of the K-S distance as determined from a Monte Carlo calculation.

The p-values from each of the distributions were determined and range from 0.001 to 0.15. The authors comment that “given the *a posteriori* selection of the 9 kinematical variables, the combined statistical significance cannot be unequivocally quantified.” However, we can determine a combined p-value by calculating the product of the 9 p-values,  $p_{tot}$ , and determining how likely it would be to obtain this product value assuming the background hypothesis. This is given by

$$P_{tot} = \prod_{m=1}^9 \left[ \sum_{k=0}^{m-1} \frac{-(\ln p_{tot})^k}{k!} \right], \quad (13)$$

and equals  $1.6 \times 10^{-6}$  assuming you set aside the reservations of the authors.

This estimate of the overall p-value raises a number of comments. First, are the variables sufficiently uncorrelated that any residual correlations can be ignored? Various tests were made of this assumption by the authors, but no rigorous argument was presented. Second, uncertainties in the Standard Model predictions have not been incorporated into the p-value calculation. These may have some effect on the overall result, but it is unclear how large this might be. Third, the effect on the p-value estimate of the *a posteriori* choice of variables is virtually impossible to assess. A study of a series of alternate variables were made by the authors, but no firm conclusion could be drawn.

Of these, perhaps the third is the most vexing. It is true that the choice of the 9 variables for this analysis was made after the 13 event data sample had been identified as being unusual. In that sense, it is no longer possible to argue that the quoted p-value is an unbiased measure of the significance of the observation.

In this case, the best strategy is to repeat the measurement with an independent data sample to determine if the same effect is observed. However, this analysis serves as a good example of the issues one must face in making such a multi-variate estimate of significance.

## 7. Observations on Current Practice and Summary

Particle physicists have increasingly relied on numerical estimates of statistical significance. The literature is replete with the use of the p-value, and this appears to have developed into one common measure, as illustrated by the examples provided above. Other measures of significance are often quoted, such as the equivalent number of standard deviations a measurement lies from the value predicted by a hypothesis. This is, of course, just a p-value under a different name.

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<sup>3</sup>The authors chose background distributions for these figures obtained using Monte Carlo calculations, but used background distributions for their p-value calculations obtained by “bootstrapping,” using a complementary data sample that had no signal events and that was argued to provide a good characterization of the expected Standard Model backgrounds. The Standard Model Monte Carlo calculations resulted in similar p-value estimates.

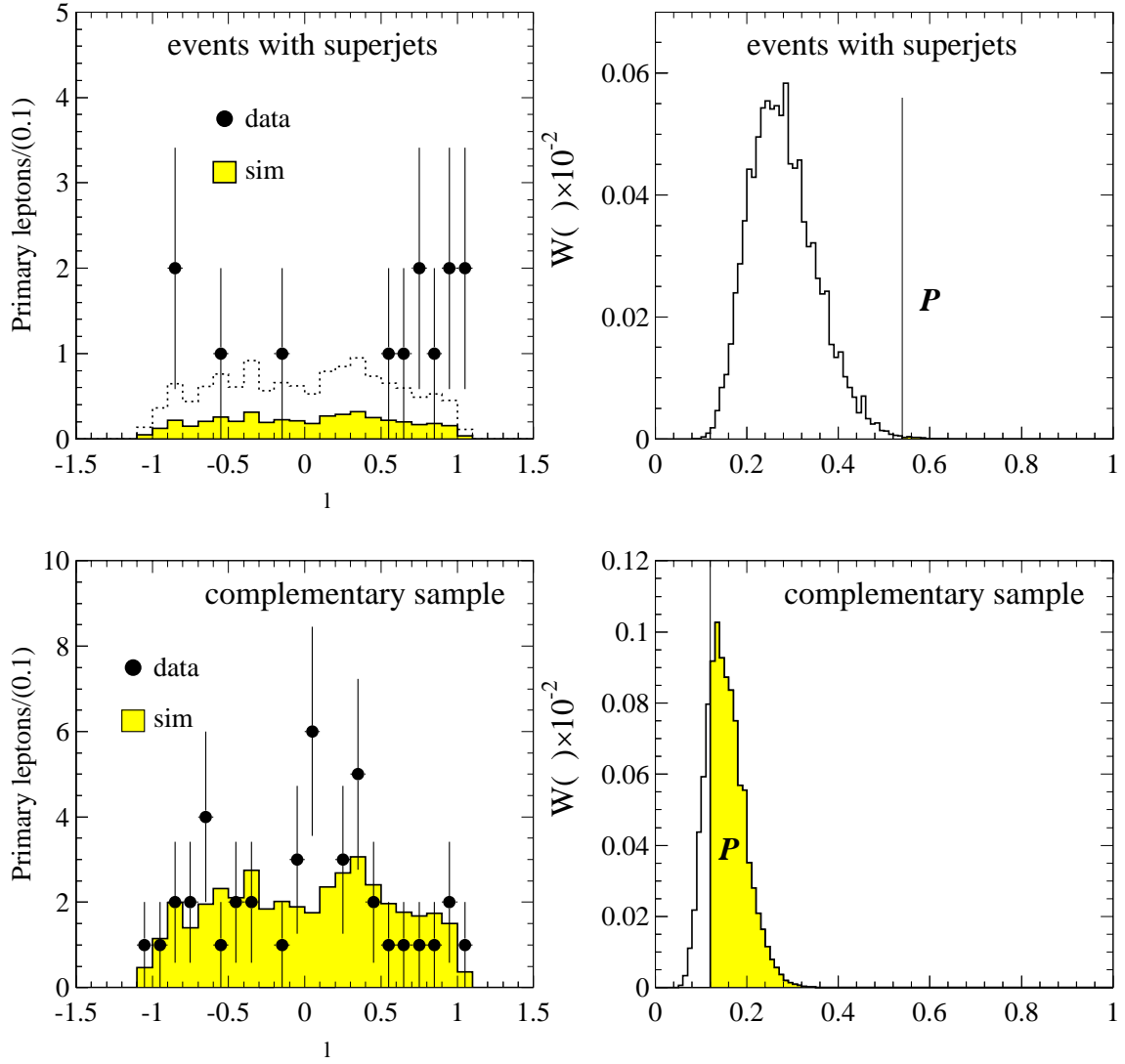


Fig. 2: The  $\eta$  distribution of the lepton from the  $W$  boson decay in the CDF “superjet” events is shown in the top-left plot (points) and compared with the Standard Model prediction (shaded histogram). The top-right distribution is the expected distribution of K-S distance of the 13 data events and the SM prediction in the top-left plot. The vertical line is the K-S distance for the two distributions. Similarly, the bottom-left plot is the lepton  $\eta$  distribution for the complementary sample of data events where a “superjet” is not detected, and the bottom-right plot gives the distribution of the corresponding K-S distance between the data and predicted distribution. The K-S test distributions were generated using a Monte Carlo calculation.

More significantly, there are consistent attempts in the literature to include in p-value estimates more complete information about a given measurement, such as the sensitivity of the estimate to systematic uncertainties and information from several statistics. The more difficult problem of avoiding unconscious bias in the selection of statistics is addressed through the use of “blind analyses,” but the effective application of such techniques to truly serendipitous discoveries is problematic. Here, the time-honoured technique of testing specific hypotheses developed through the study of one data set by creating and analyzing an independent data set with at least comparable statistical power remains the most effective tool for separating what we would call the “statistical fluctuations” from first evidence for truly new phenomena.

Finally, what is an appropriate criteria for claiming a discovery on the basis of the p-value of the null hypothesis? The recent literature would suggest a p-value in the range of  $10^{-6}$ , comparable to a “ $5\sigma$ ” observation, provides convincing evidence. However, the credibility of such a claim relies on the care taken to avoid unconscious bias in the selection of the data and the techniques chosen to calculate the p-value.

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